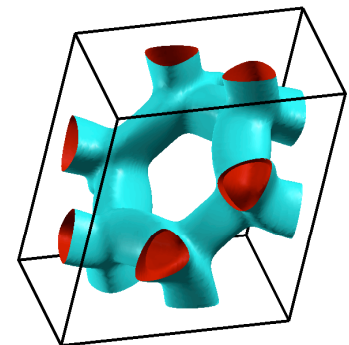
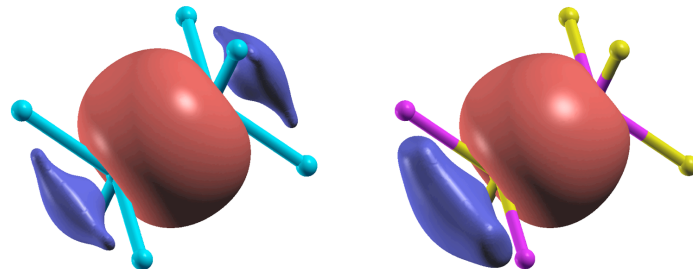
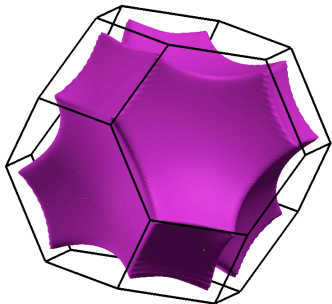


# Wannier Functions: *ab-initio* tight-binding

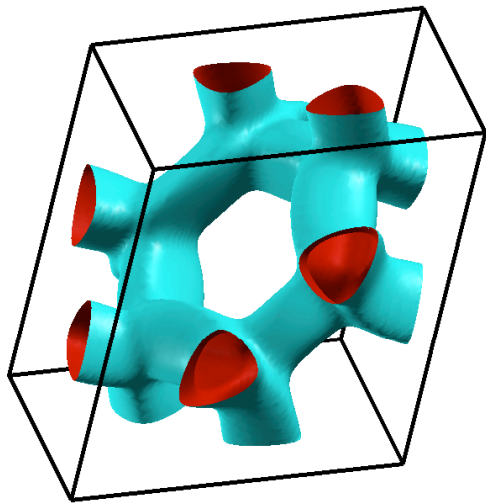
Jonathan Yates

*Cavendish Laboratory, Cambridge University*

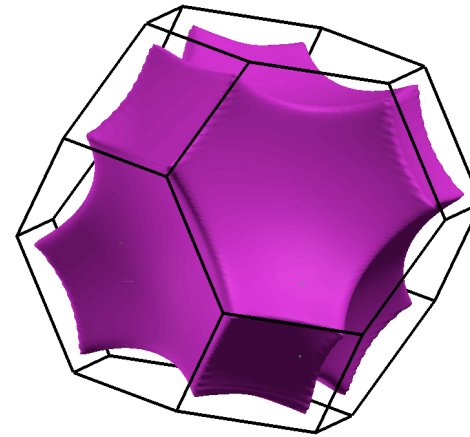


# Representing a Fermi Surface

Accurate description of Fermi surface properties requires a detailed sampling of the Brillouin Zone



Lead Fermi surface



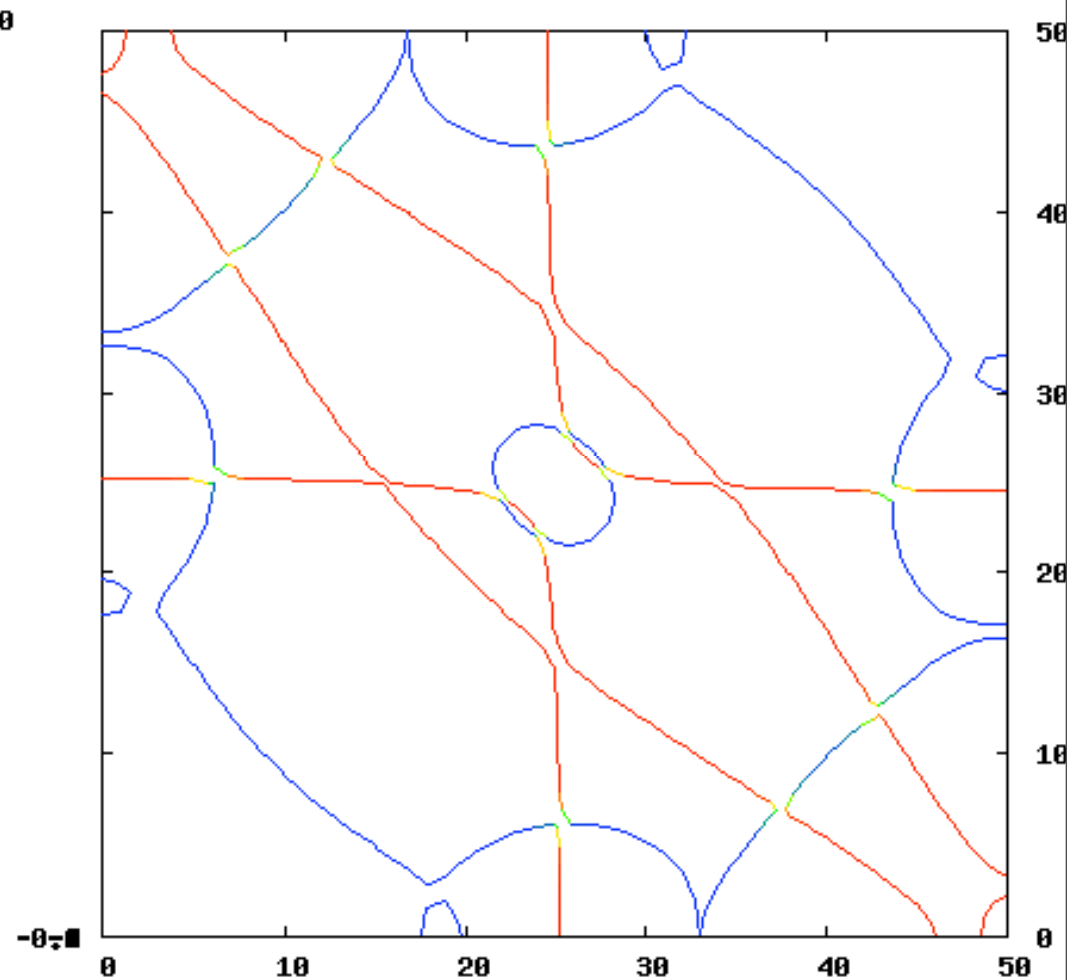
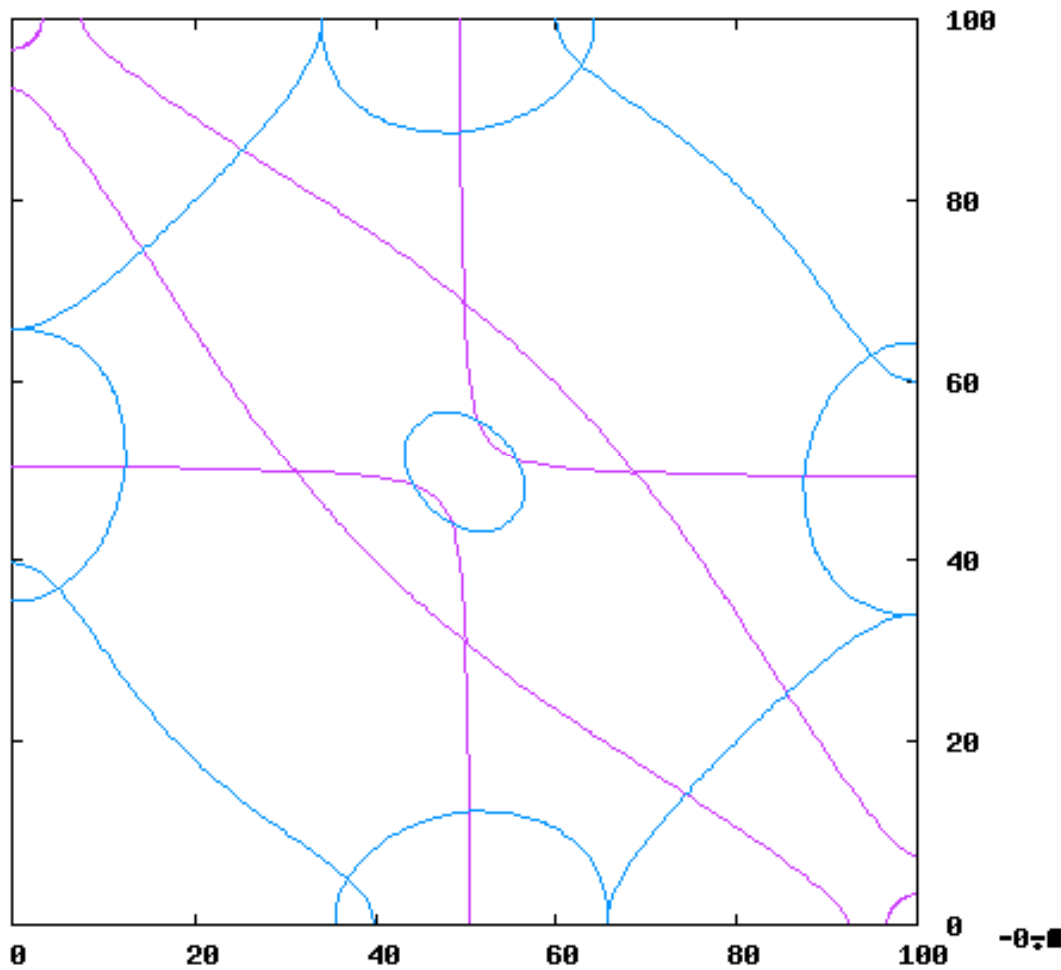
Al Fermi surface

# Representing a Fermi Surface

Fe fermi surface

Scalar Relativistic

In the  $(010)$  plane with spin-orbit coupling



# Outline

- **Wannier Functions**

- *One band*
- *Isolated Set of Bands*
- *Entangled Bands*

- **Wannier Interpolation**

- *Accurate and Efficient approach to Fermi surface and spectral properties*

- **Examples**

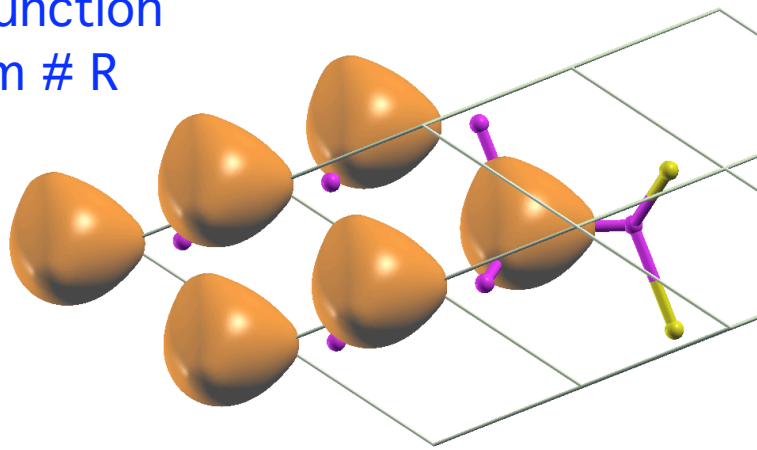
- *Anomalous Hall Effect*
- *Electron-Phonon Coupling*

# Wannier Functions

$$w_{n\mathbf{R}}(\mathbf{r}) = \frac{V}{(2\pi)^3} \int_{BZ} \psi_{n\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$

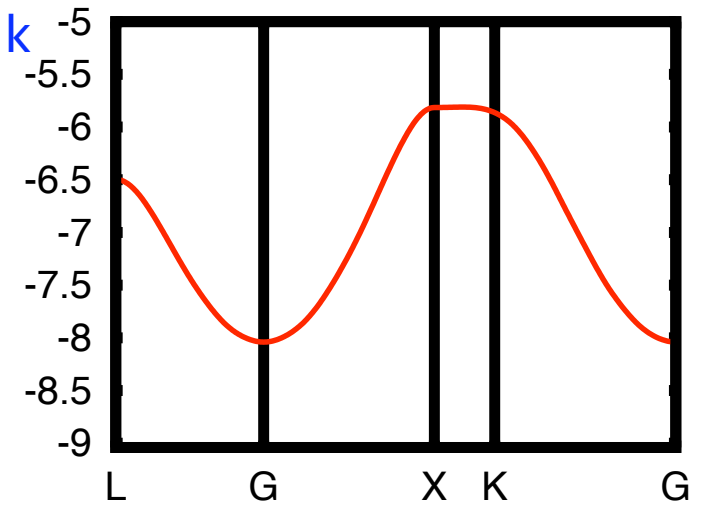
↑  
Wannier function  
Quantum # R

← Bloch state  
Quantum # k



$$\psi_{n\mathbf{k}}(\mathbf{r}) \rightarrow e^{-i\phi(\mathbf{k})} \psi_{n\mathbf{k}}(\mathbf{r})$$

GaAs

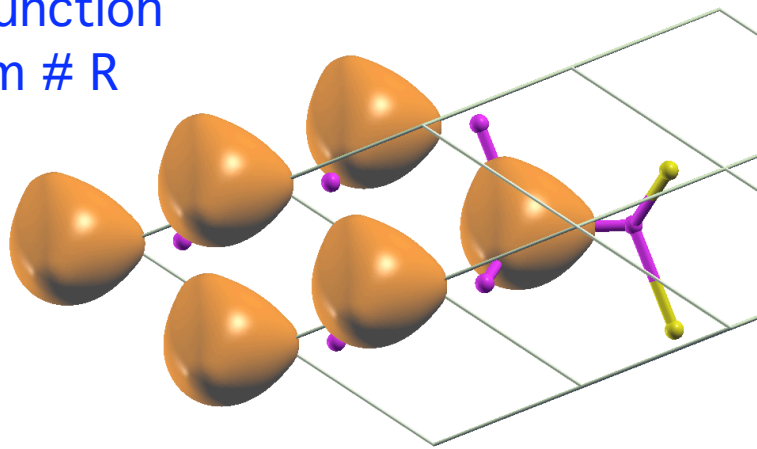


# Wannier Functions

$$w_{n\mathbf{R}}(\mathbf{r}) = \frac{V}{(2\pi)^3} \int_{BZ} \psi_{n\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$

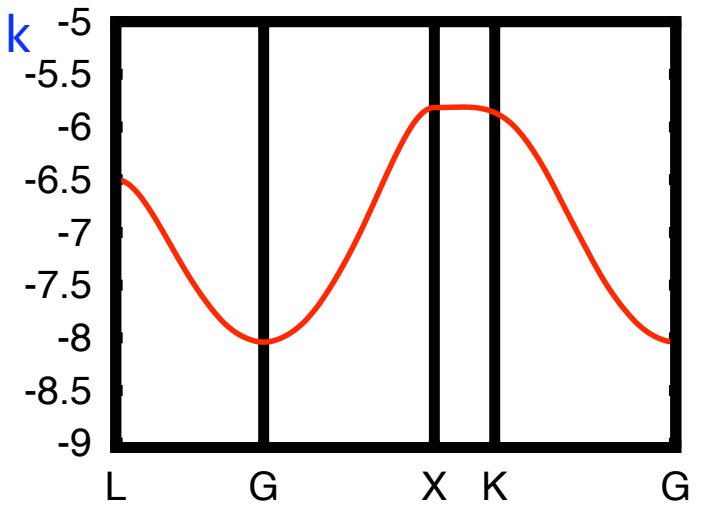
↑  
Wannier function  
Quantum # R

← Bloch state  
Quantum # k



$$\psi_{n\mathbf{k}}(\mathbf{r}) \rightarrow e^{-i\phi(\mathbf{k})} \psi_{n\mathbf{k}}(\mathbf{r})$$

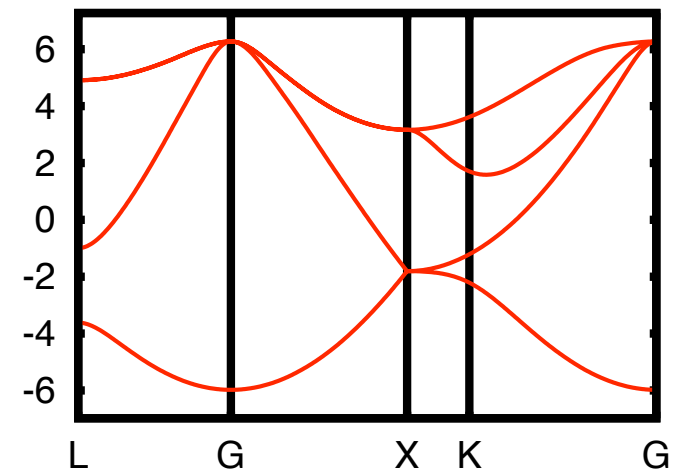
GaAs



## Multiband - Generalized WF

$$\psi_{n\mathbf{k}}(\mathbf{r}) \rightarrow \sum_n U_{mn}^{(\mathbf{k})} \psi_{m\mathbf{k}}(\mathbf{r})$$

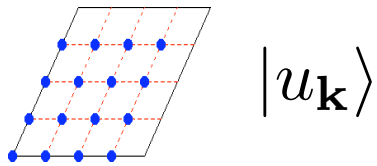
Si



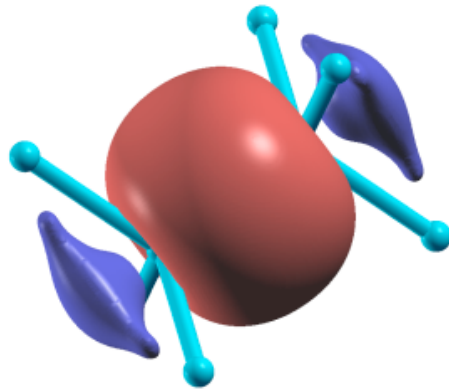
# Maximally Localised Wannier Functions

$$w_{n\mathbf{R}}(\mathbf{r}) = \frac{V}{(2\pi)^3} \int_{BZ} \left[ \sum_m U_{mn}^{(\mathbf{k})} \psi_{m\mathbf{k}}(\mathbf{r}) \right] e^{-\mathbf{k} \cdot \mathbf{R}} d\mathbf{k}$$

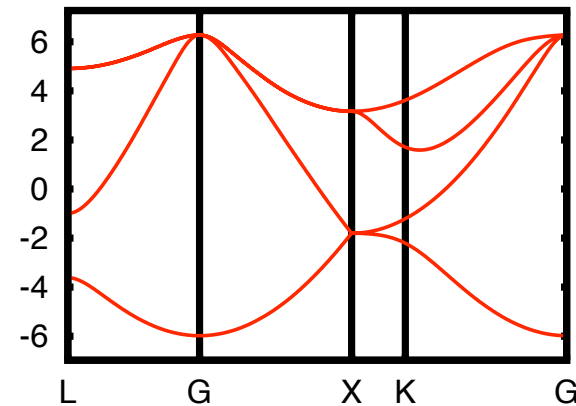
Choose  $U^{\mathbf{k}}$  to minimise quadratic spread (Marzari, Vanderbilt)



WF defined in basis of Bloch states



Si valence bands

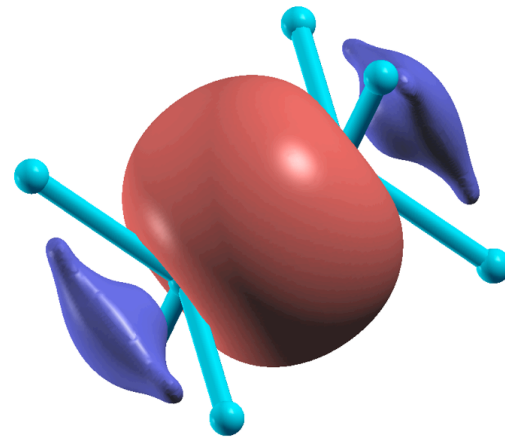


Wannier **localisation** in  $\mathbf{R}$  gives Bloch **smoothness** in  $\mathbf{k}$

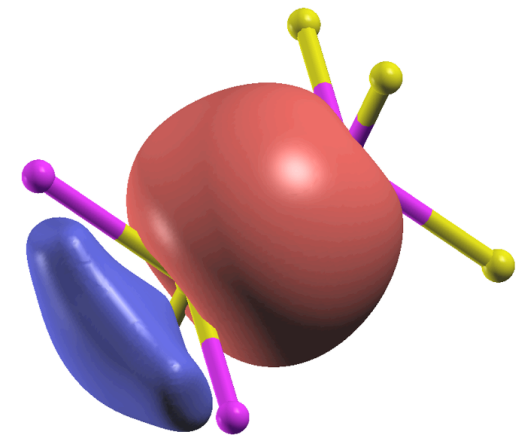
$$\psi_{n\mathbf{k}}^{rot}(\mathbf{r}) = \sum_m U_{mn}^{(\mathbf{k})} \psi_{m\mathbf{k}}(\mathbf{r})$$

# Wannier Functions: A local picture

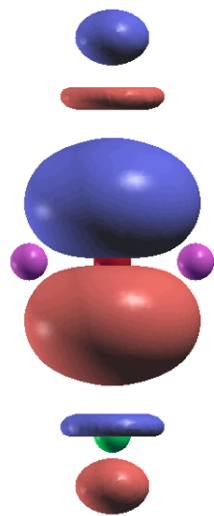
1- Intuitive picture of local bonding



Si

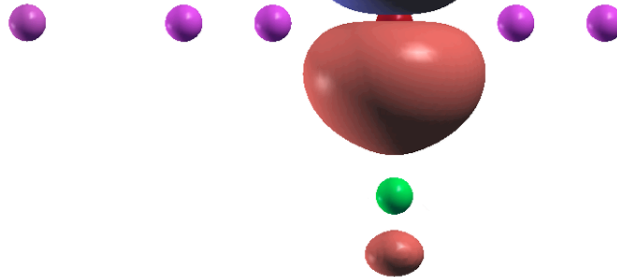


GaAs



Paraelectric

BaTiO<sub>3</sub>



Ferroelectric

2- Wannier centres give bulk polarisation in a ferroelectric

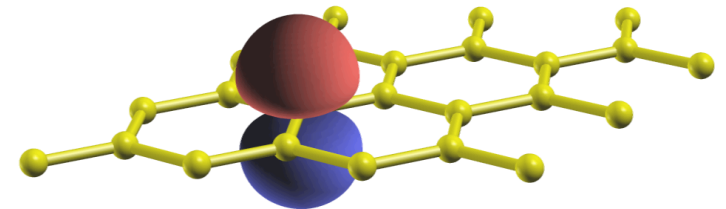
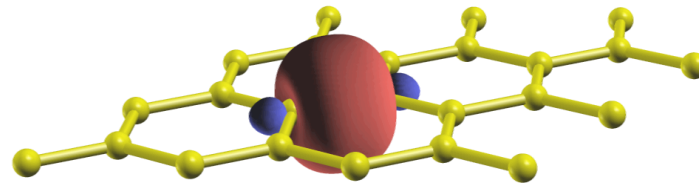
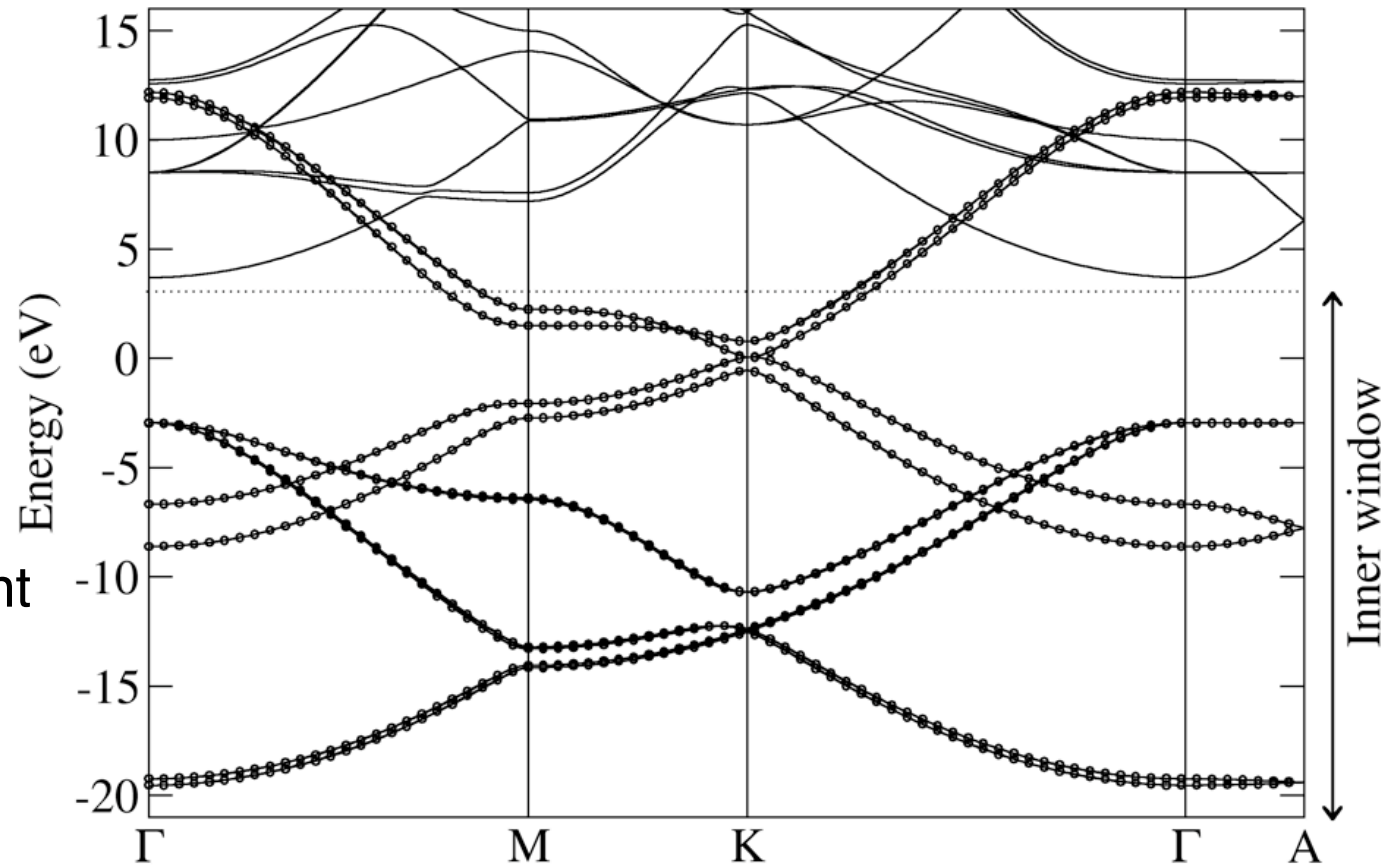


# MLWF - Entangled Bands

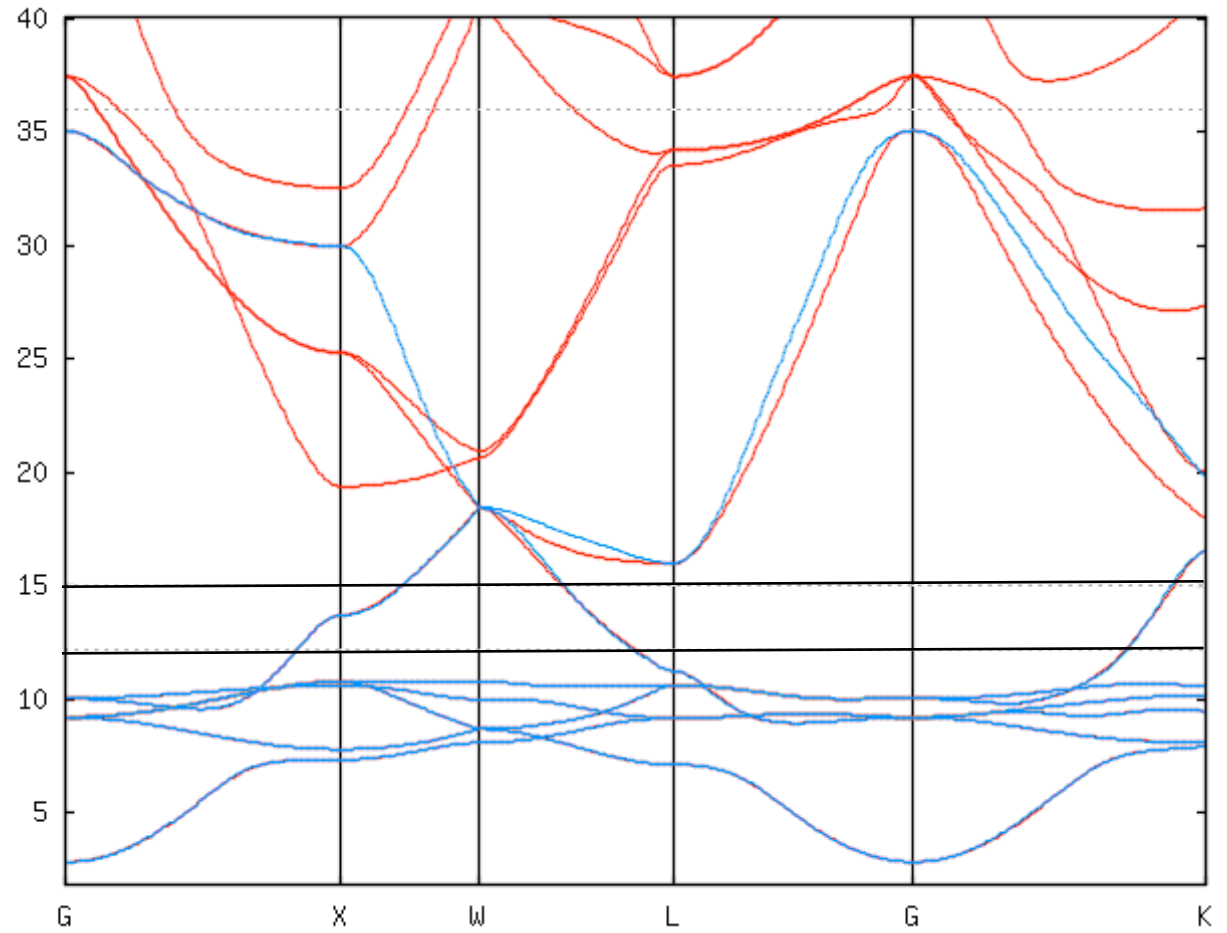
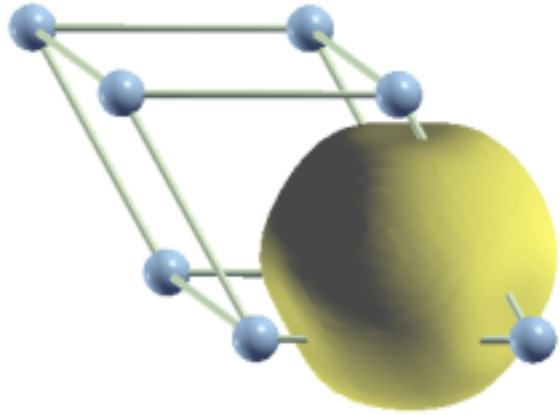
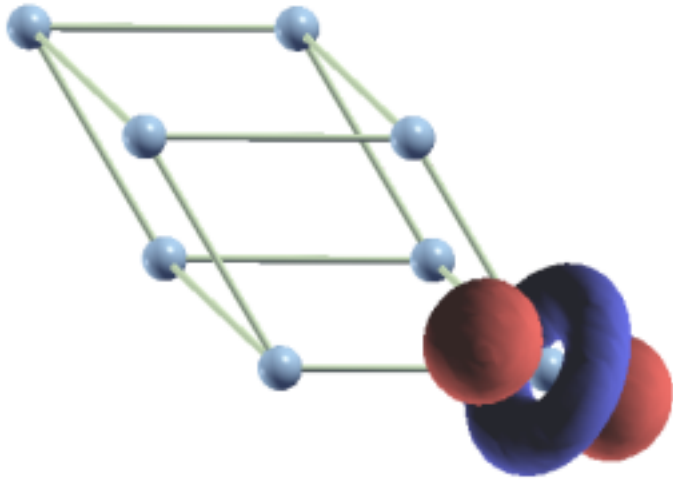
Disentanglement procedure  
(Souza, Marzari and Vanderbilt)

Obtain “partially occupied”  
Wannier functions

Essentially perfect agreement  
within given energy window



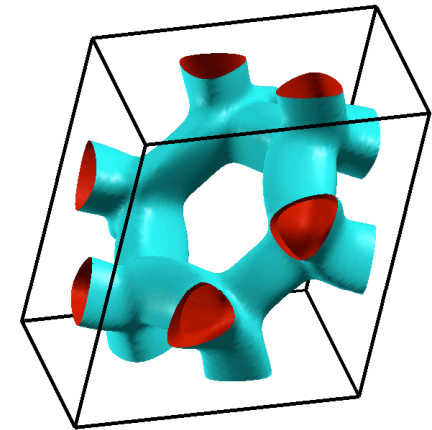
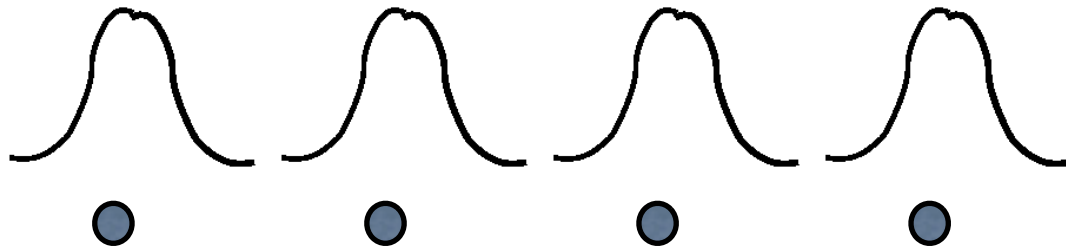
# Copper



# Wannier Interpolation

*Accurate and Efficient approach to Fermi surface and spectral properties*

- Exploit **localisation** of Wannier functions.
- Require only matrix elements between close neighbours.



Lead Fermi surface

1st principles accuracy at tight-binding cost  
Interpolation of any one-electron operator

# Wannier Interpolation

1- Obtain operator in Wannier basis

$$O_{mn}(\mathbf{R}) = \langle Om | \hat{O} | Rn \rangle$$

# Wannier Interpolation

1- Obtain operator in Wannier basis

$$O_{mn}(\mathbf{R}) = \langle Om | \hat{O} | Rn \rangle$$

2- Fourier transform to an arbitrary k-point,  $\mathbf{k}'$

$$O_{mn}(\mathbf{k}') = \sum_{\mathbf{R}} e^{i\mathbf{k}' \cdot \mathbf{R}} O_{mn}(\mathbf{R})$$

# Wannier Interpolation

1- Obtain operator in Wannier basis

$$O_{mn}(\mathbf{R}) = \langle Om | \hat{O} | Rn \rangle$$

2- Fourier transform to an arbitrary k-point,  $\mathbf{k}'$

$$O_{mn}(\mathbf{k}') = \sum_{\mathbf{R}} e^{i\mathbf{k}' \cdot \mathbf{R}} O_{mn}(\mathbf{R})$$

3- Un-rotate matrix at  $\mathbf{k}'$

$$\langle \psi_{m,\mathbf{k}'} | O | \psi_{n,\mathbf{k}'} \rangle = [U_{\mathbf{k}'}^\dagger, O(\mathbf{k}') U_{\mathbf{k}'}]$$

What we need

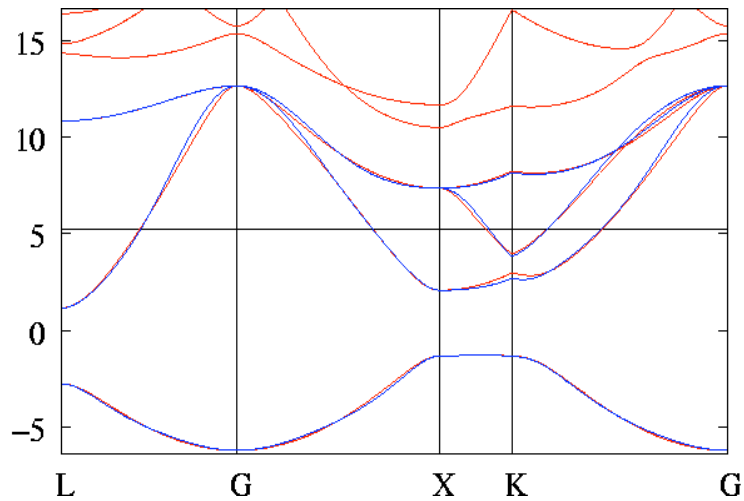
diagonalises H

All operations involve small matrices ( $N_{\text{wan}} \times N_{\text{wan}}$ )

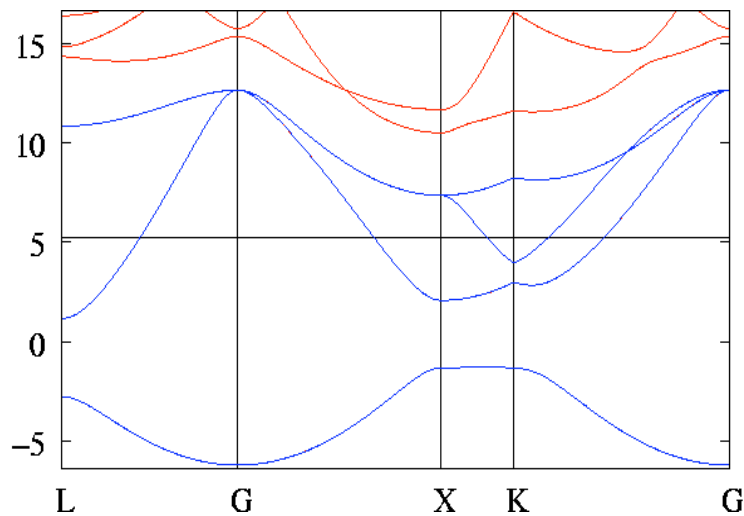
=> FAST!

# Convergence

Lead



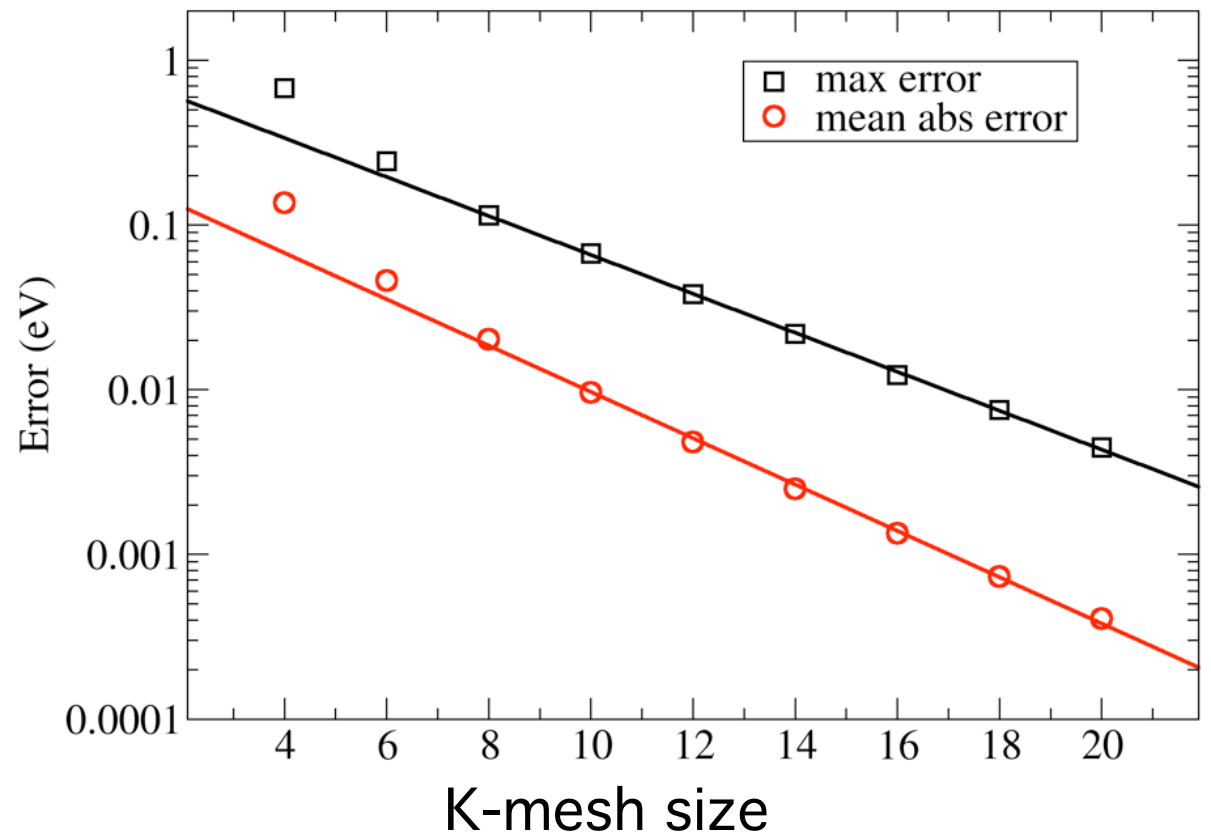
4x4x4



10x10x10

Wannier interpolated object converges to ab-initio value, **exponentially** with k-grid sampling.

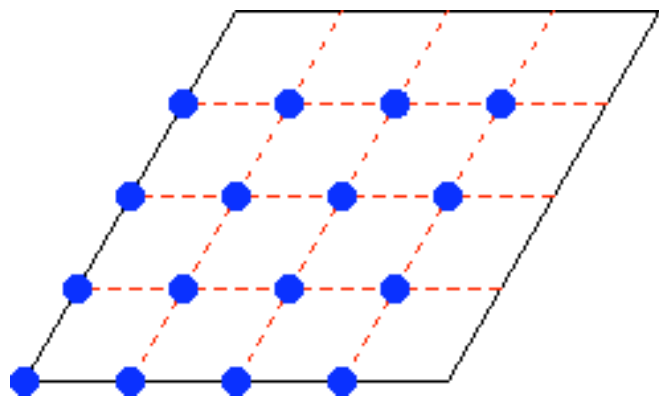
Band interpolation error



# Wannier Interpolation - Summary

## Bloch States eg plane wave basis

- Each k-point expensive
- Calculation of MLWF converges exponentially

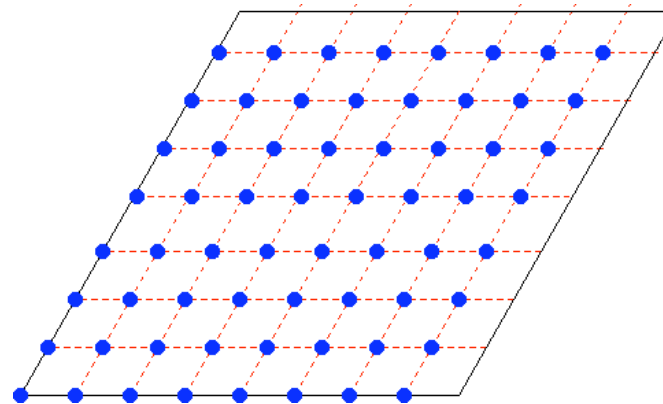


Sample on coarse grid



## Wannier basis

- Each k-point cheap
- Introduce occupancies and compute properties (only polynomial convergence)

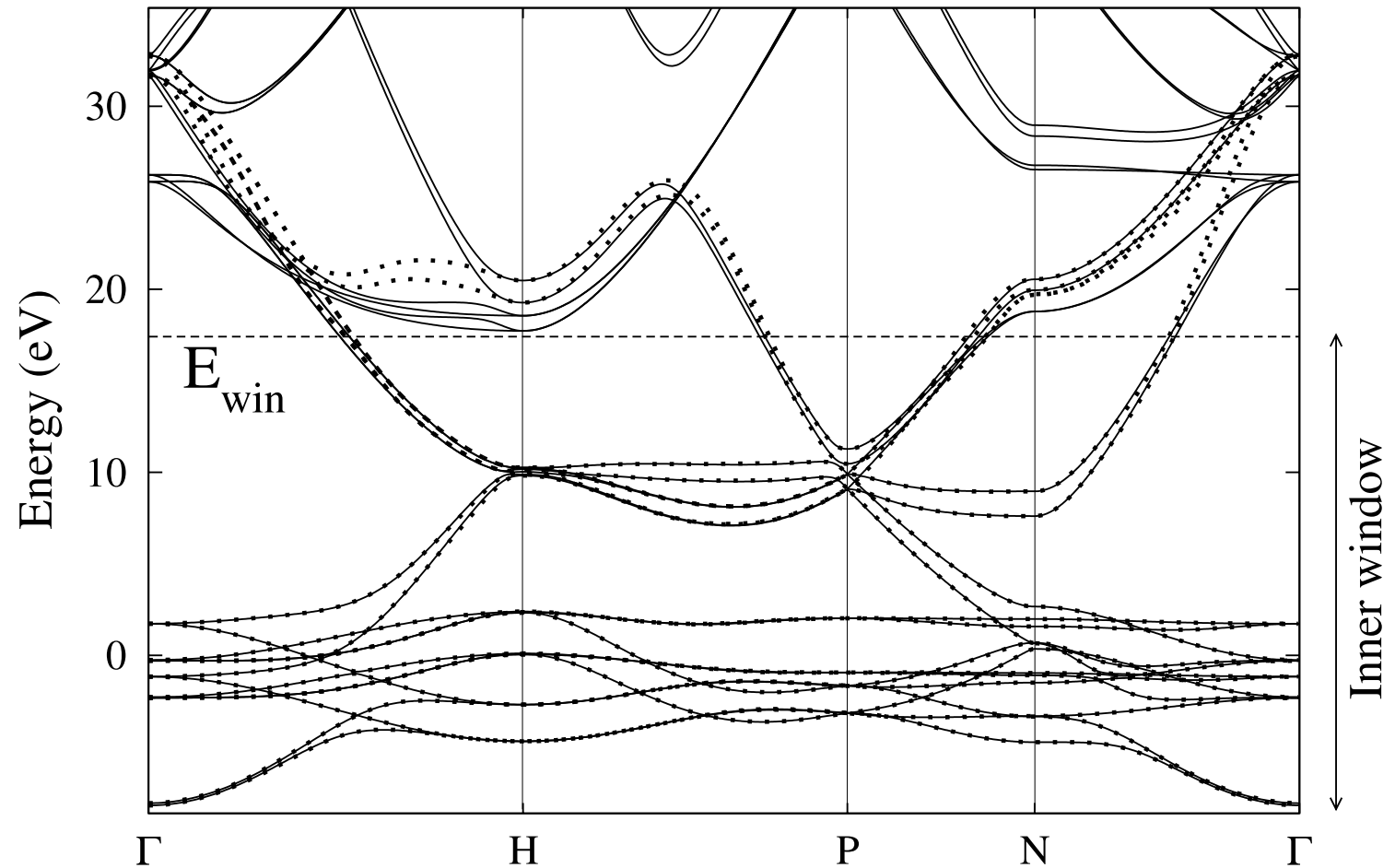
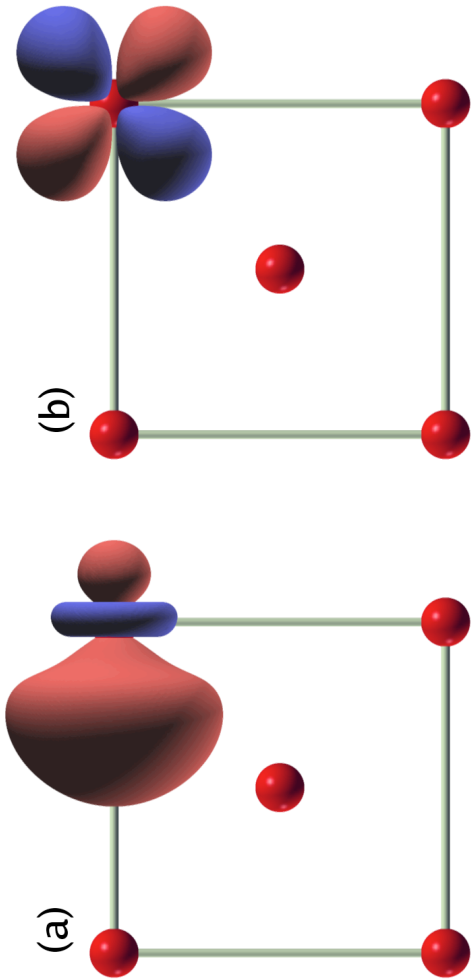


Sample on fine grid

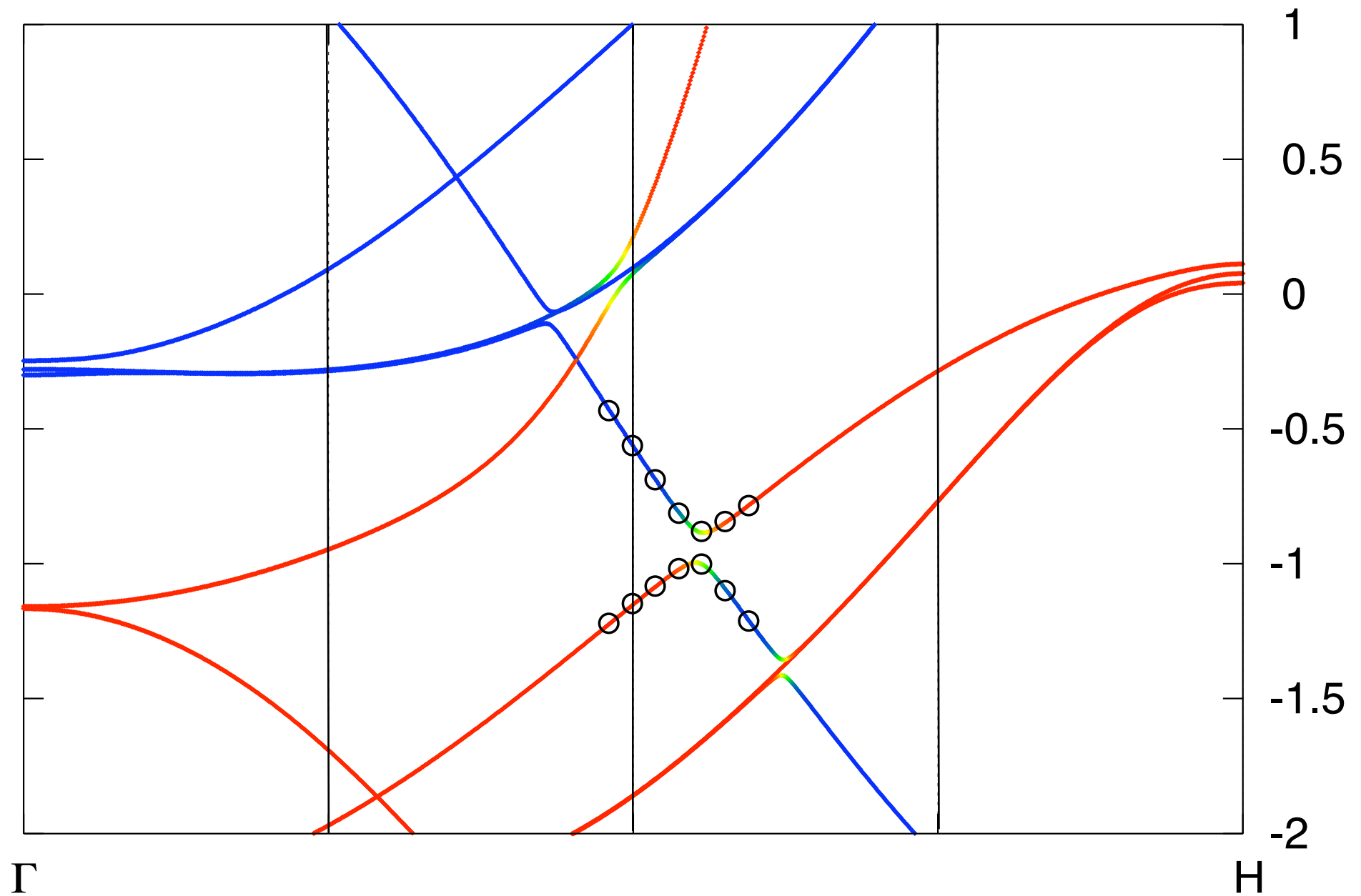


# Wannier Interpolation - Fe

- 18 spinor Wannier functions
- Keep up to 4th neighbour overlaps
- Cost 1/2000 of full calculation



# Wannier Interpolation



# Wannier Interpolation pt2

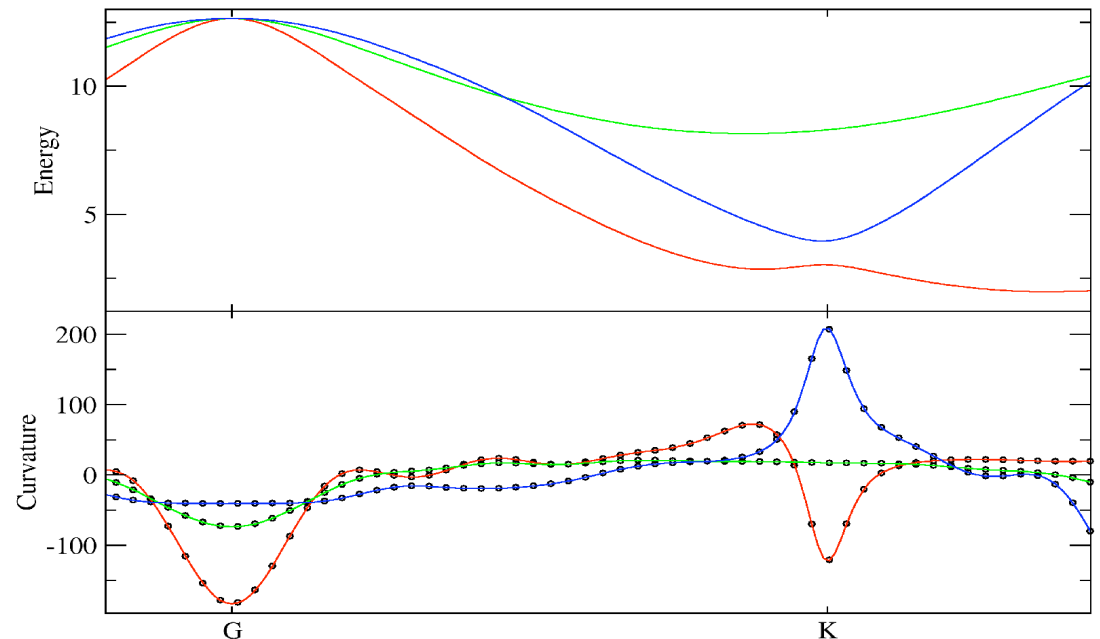
1- k-derivatives can be taken analytically

eg band gradient

$$H_{\mathbf{k}} = \sum_{\mathbf{R}} e^{i\mathbf{k}\cdot\mathbf{R}} H(\mathbf{R})$$

$$\frac{\partial H_{\mathbf{k}}}{\partial \mathbf{k}_{\alpha}} = i \sum_{\mathbf{R}} R_{\alpha} e^{i\mathbf{k}\cdot\mathbf{R}} H(\mathbf{R})$$

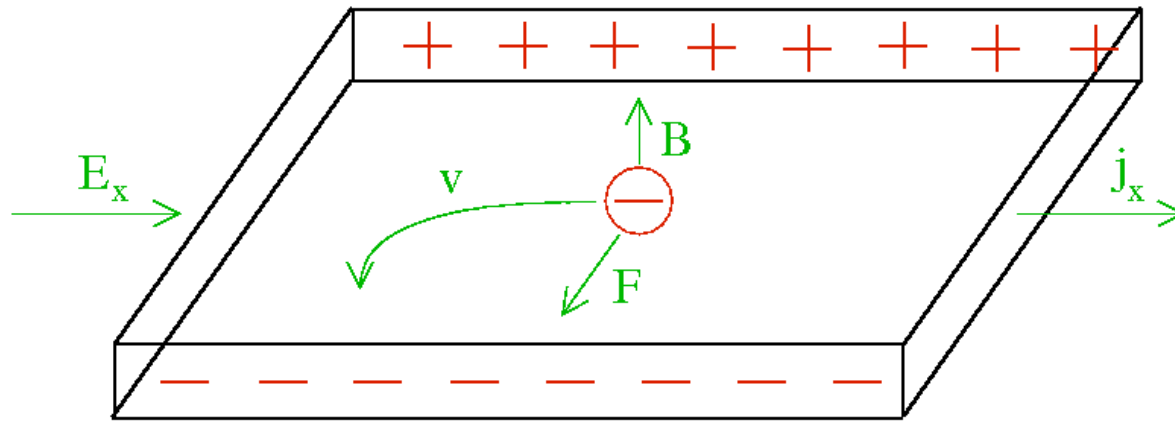
repeat for higher derivatives



2- Position operator matrix elements well defined

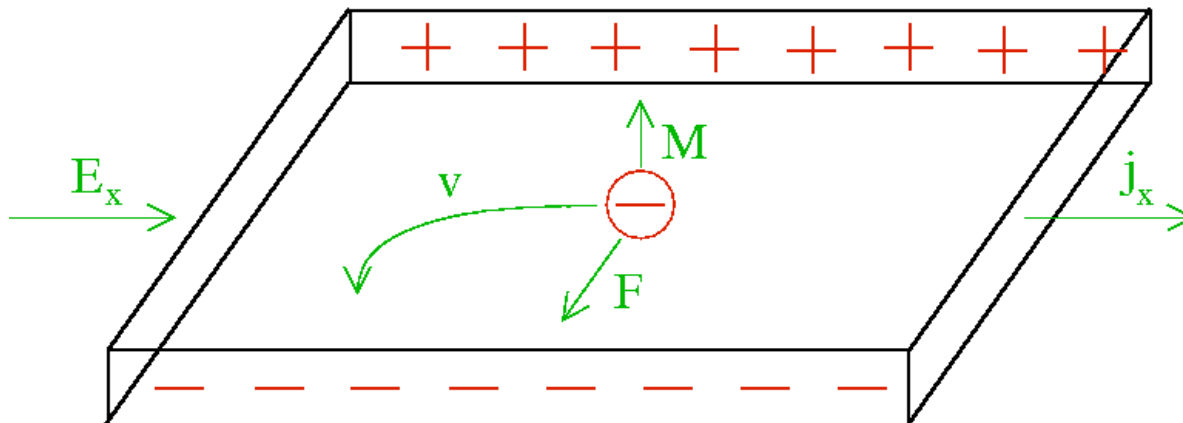
# Hall Effect

## Ordinary Hall Effect (Hall 1879)



Crossed E and B fields  
B breaks time-reversal

## Anomalous Hall Effect (Hall 1880)



No B field needed!  
Ferromagnetism  
breaks time-reversal

# Mechanisms for the AHE

## Intrinsic

1954 Karplus and Luttinger

property of electron motion in perfect lattice

“dissipationless” current (spin-orbit effect)

## Extrinsic

1955 Smit “skew scattering”

1970 Berger “side-jump”

## Intrinsic (again)

1996- Reexamination of KL in terms of Berry’s phases

Niu (Texas), Nagaosa (Tokyo)

# Electron Dynamics

Textbook wavepacket dynamics

$$\dot{\mathbf{r}} = \frac{1}{\hbar} \frac{\partial \varepsilon_{n\mathbf{k}}}{\partial \mathbf{k}}$$

$$\hbar \dot{\mathbf{k}} = -e\mathbf{E} - e\dot{\mathbf{r}} \times \mathbf{B}$$

# Electron Dynamics

Textbook wavepacket dynamics missing a term: “anomalous velocity”

$$\dot{\mathbf{r}} = \frac{1}{\hbar} \frac{\partial \varepsilon_{n\mathbf{k}}}{\partial \mathbf{k}} - \dot{\mathbf{k}} \times \boldsymbol{\Omega}_n(\mathbf{k})$$

“k-space magnetic field” caused by lattice potential

“k-space Lorentz force” even when B=0

$$\hbar \dot{\mathbf{k}} = -e\mathbf{E} - e\dot{\mathbf{r}} \times \mathbf{B}$$

## Anomalous Hall Conductivity

$$\sigma_{xy} = \frac{-e^2}{(2\pi)^2 \hbar} \sum_n \int_{\text{BZ}} d\mathbf{k} f_n(\mathbf{k}) \Omega_{n,z}(\mathbf{k})$$

Fermi function

Berry Curvature

# Berry Curvature

Berry  
Curvature

$$\Omega_n(\mathbf{k}) = -\text{Im} \langle \nabla_{\mathbf{k}} u_{n,\mathbf{k}} | \times | \nabla_{\mathbf{k}} u_{n,\mathbf{k}} \rangle$$

cell periodic  
Bloch state

Berry  
Potential

$$\mathbf{A}_n(\mathbf{k}) = i \langle u_{n\mathbf{k}} | \nabla_{\mathbf{k}} | u_{n\mathbf{k}} \rangle \quad \Omega_n(\mathbf{k}) = \nabla \times \mathbf{A}_n(\mathbf{k})$$

## Ab-initio Calculation

Finite-differences difficult: phases, band-crossings

**bcc iron** (LAPW) Yao et al (PRL 2004)

Kubo formula

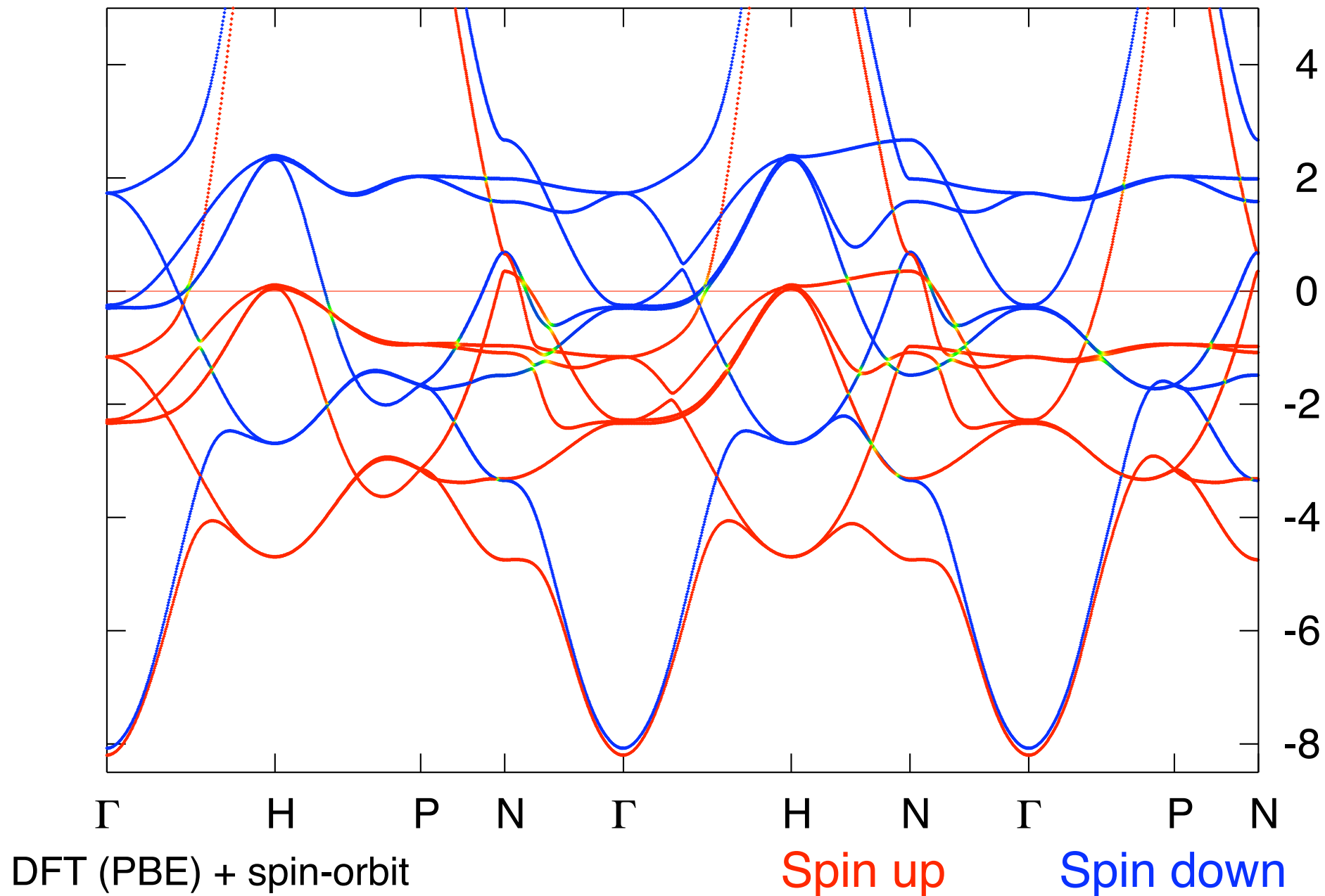
$$\Omega_{n,z}(\mathbf{k}) = -2\text{Im} \sum_{m \neq n} \frac{v_{nm,x}(\mathbf{k}) v_{mn,y}(\mathbf{k})}{(\omega_m(\mathbf{k}) - \omega_n(\mathbf{k}))^2}$$

Extremely computationally expensive!

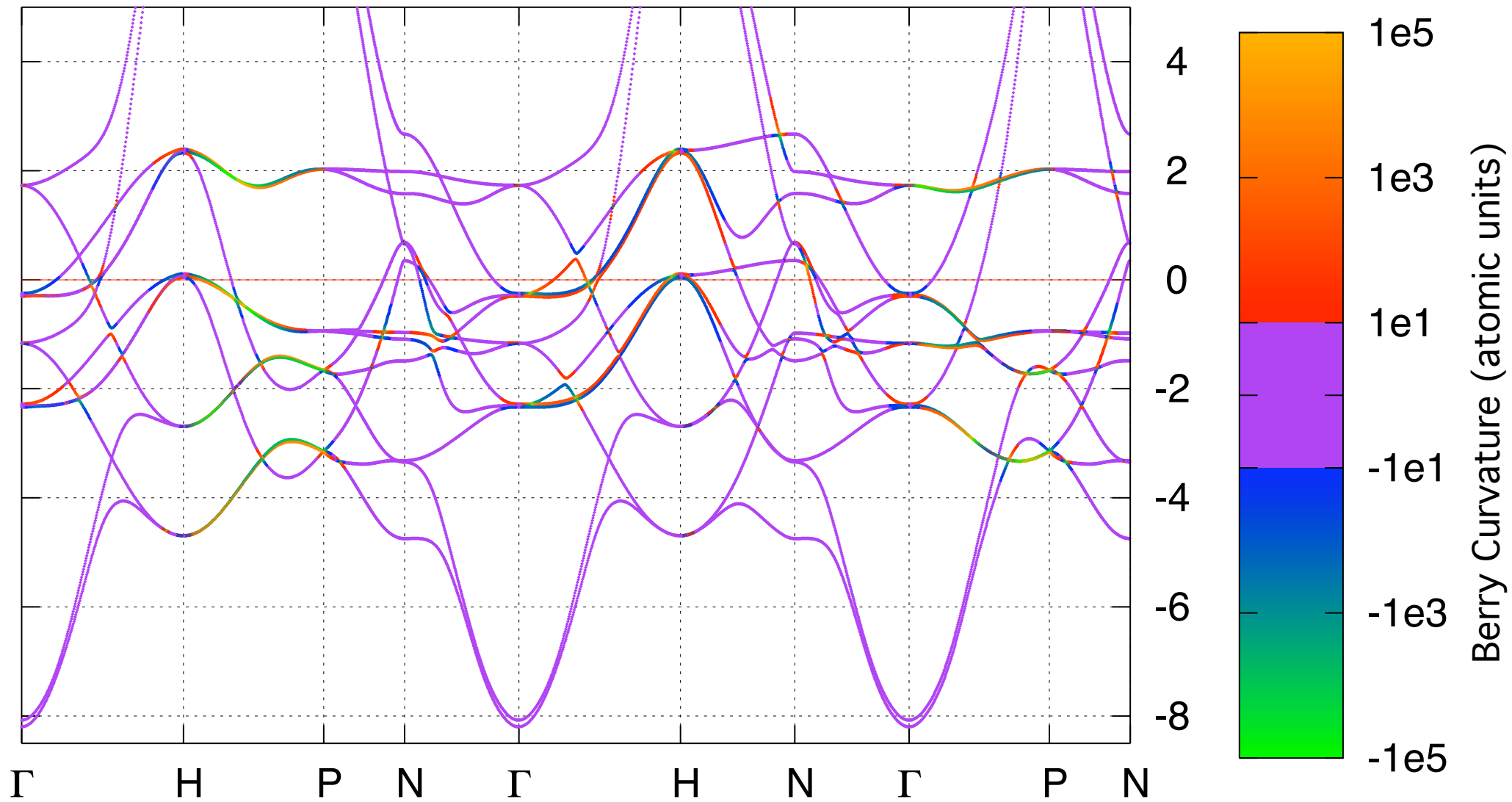
final number within 20% of experiment



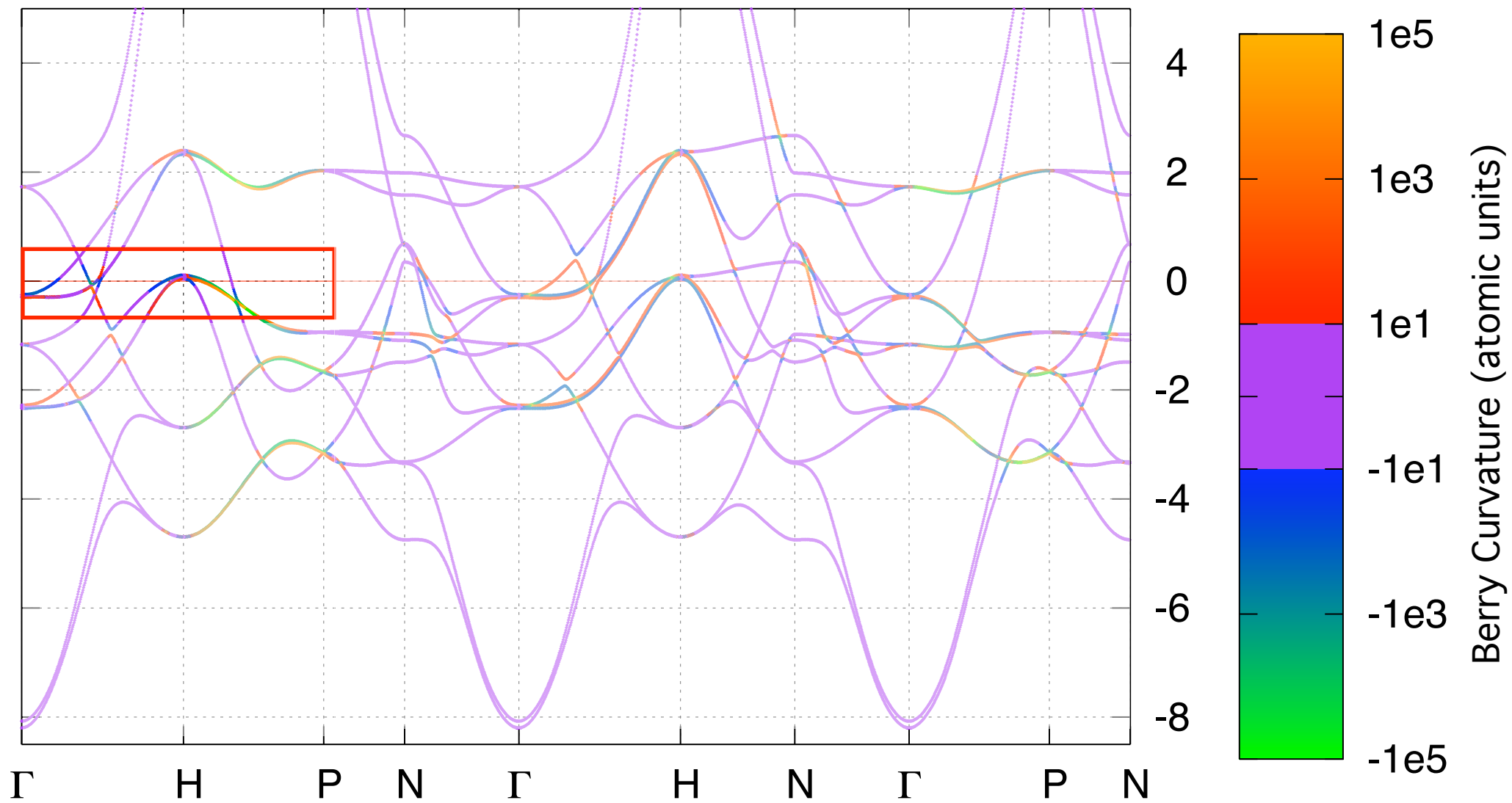
# Band-structure bcc Fe



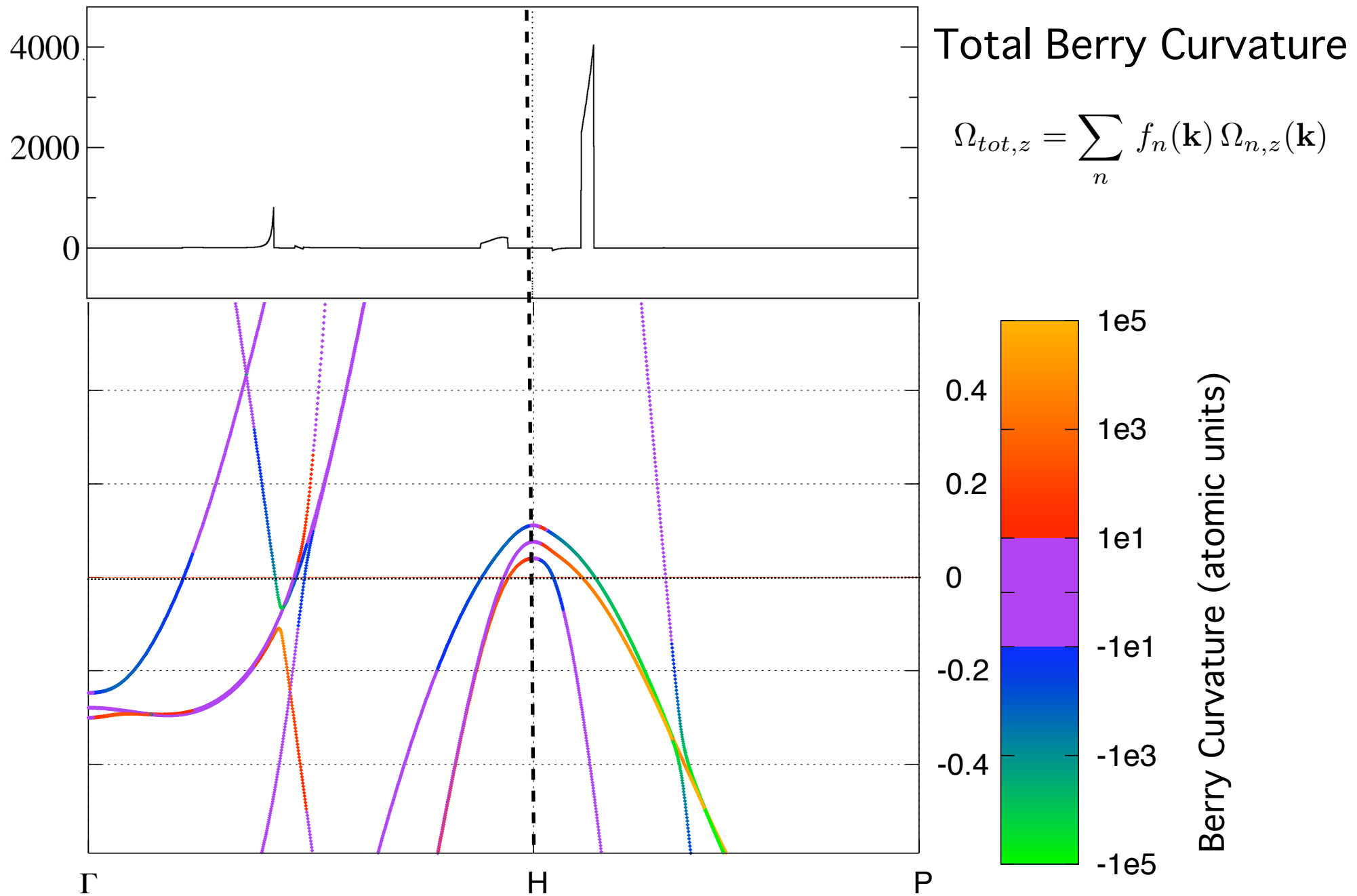
# Berry curvature bcc Fe



# Berry curvature bcc Fe



# Berry Curvature



# Our Calculation

Ab-initio k-grid: 8x8x8

Interpolation k-grid: 320x320x320 + 7x7x7 “adaptive refinement”  
(when  $\Omega > 100$  a.u.)

## Wannier Interpolation

planewaves +  
pseudopotentials

758  $(\Omega \text{ cm})^{-1}$

## Existing Calculation\*

Kubo formula +  
LAPW

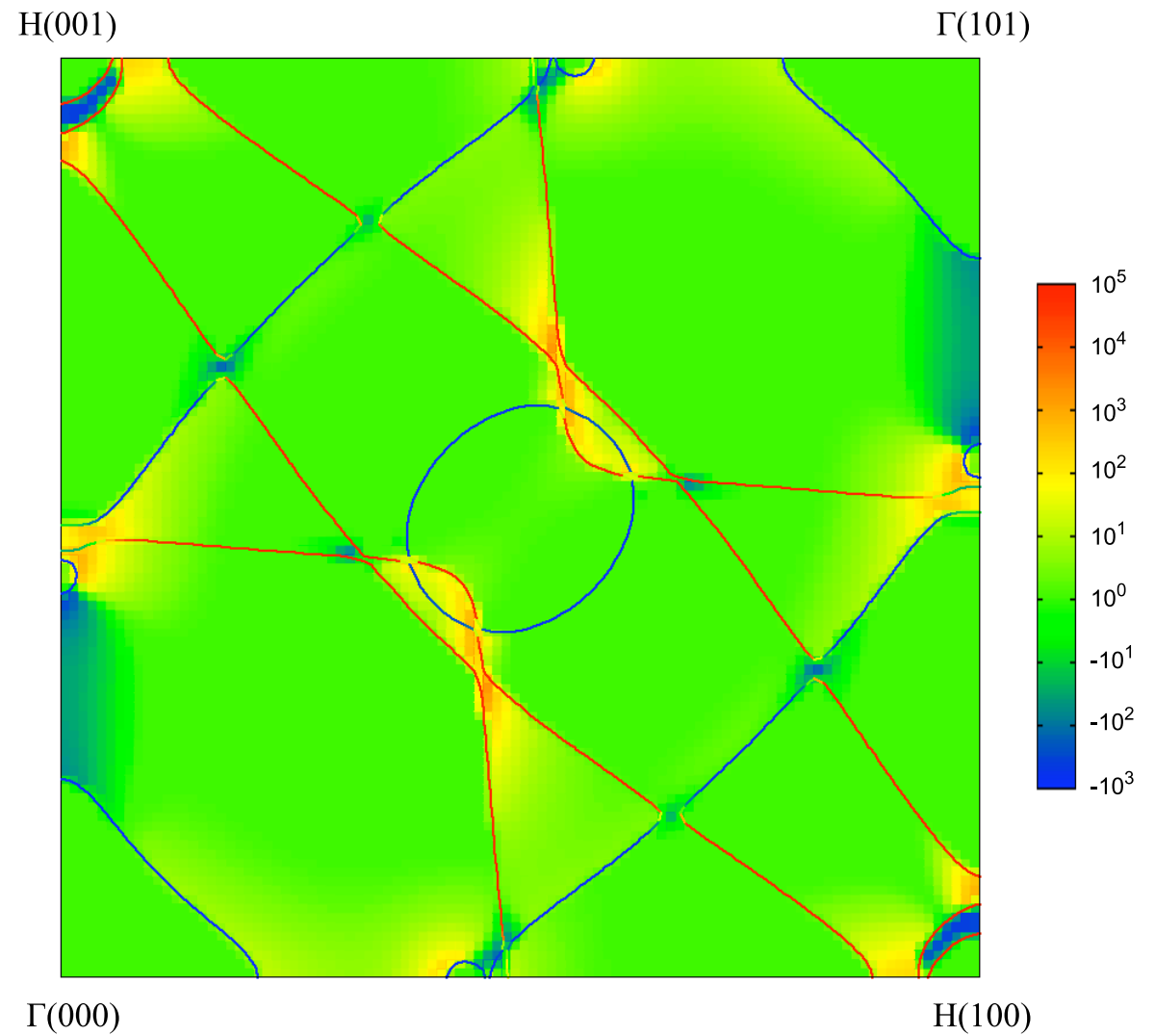
751  $(\Omega \text{ cm})^{-1}$

Experiment  $\sim 1000 (\Omega \text{ cm})^{-1}$

Differences: Scattering, Approximate DFT, Experimental uncertainty

\*Yao et al PRL (2004)

# Berry curvature in $k_y=0$ plane

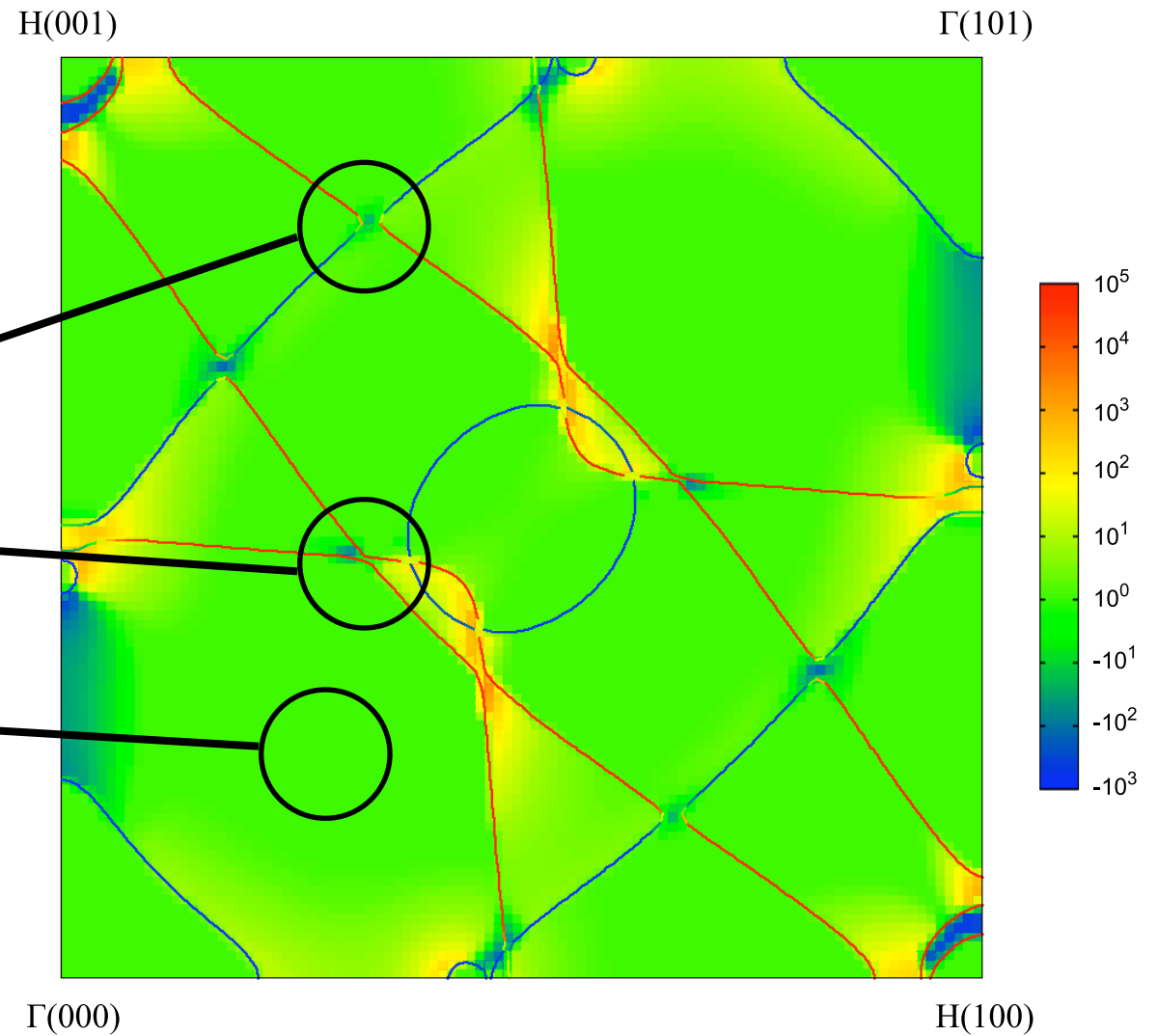


Total Berry curvature (atomic units)

# Berry curvature in $k_y=0$ plane

## Contributions

26%	opposite spin
21%	like spin
53%	smooth background



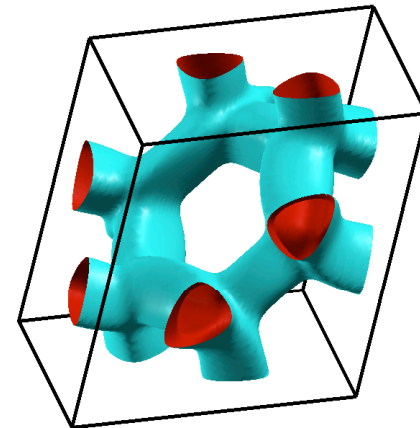
Total Berry curvature (atomic units)

# Wannier Interpolation

- Map low energy electronic structure onto “exact tight binding model”
- Compute quantities with ab-initio accuracy and tight-binding cost

Basis set independent (planewaves, LAPW, guassians etc)  
Any single particle level of theory (LDA, LDA+U, GW)

- Many applications:
  - Accurate DOS, joint-DOS, Fermi Surfaces
  - Spin-relaxation
  - Magneto-crystalline anisotropy
  - NMR in metals (Knight shift)
  - Electron-phonon interaction





# Acknowledgments

**Ivo Souza** : UC Berkeley

Wannier Interpolation: Phys Rev B 75 195121 (2007)

Anomalous Hall Effect

**Xinjie Wang, David Vanderbilt** : Rutgers University

Phys Rev B 74 195118 (2006)

Wannier90 code

**Arash Mostofi, Nicola Marzari** MIT

Released under the GPL at [www.wannier.org](http://www.wannier.org)

Electron-Phonon

**Feliciano Guistino, Marvin Cohen, Steven Louie** : UCB

Phys Rev Lett 94 047005 (2007)

Funding - Lawrence Berkeley National Laboratory  
Corpus Christi College, Cambridge