

**HOMEWORK SET 04**  
 Theory of Condensed Matter  
 UFV/TKL1/99 lecture by Martin Gmitra  
 Winter Semester 2021, room KNKTFA(Pa9-PKn)

1. [1 point] Show that matrix element of the momentum operator  $\mathbf{p}$  in the basis of the Bloch states is  $\mathbf{p}_{nn'} = \hbar\mathbf{k}\delta_{nn'} + \mathbf{p}_{nn'}^u$ , where  $\mathbf{p}_{nn'}^u = \langle u_n | -i\hbar\nabla | u_{n'} \rangle$ .
2. [1 point] Show that  $\langle u_{n,\mathbf{k}_0} | H' | u_{n,\mathbf{k}_0} \rangle = 0$  for  $\mathbf{k}_0 = 0$ , where the perturbation  $H' = \frac{\hbar}{m}\mathbf{k} \cdot (\mathbf{p} + \hbar\mathbf{k}_0)$  and  $u_{n,\mathbf{k}}(\mathbf{r})$  is the periodic part of the Bloch function.
3. [2 points] We know that the Schrödinger equation for a periodic potential  $U(\mathbf{r})$  can be written in the form  $H(\mathbf{k})u_{n\mathbf{k}}(\mathbf{r}) = \varepsilon_{n\mathbf{k}}u_{n\mathbf{k}}(\mathbf{r})$  where  $H(\mathbf{k}) = \frac{\hbar^2}{2m}(-i\nabla + \mathbf{k})^2 + U(\mathbf{r})$ . Suppose we know energies for the Bloch states at the  $\mathbf{k} = 0$ , that is the solution of  $H(0)u_{n0}(\mathbf{r}) = \varepsilon_{n0}u_{n0}(\mathbf{r})$ . Since  $u_{n0}(\mathbf{r})$  form a complete basis of functions periodic with lattice, we can expand  $u_{n\mathbf{k}}(\mathbf{r}) = \sum_{n'} c_{nn'}(\mathbf{k})u_{n'0}(\mathbf{r})$ . Find the eigenvalue problem for the coefficients  $c_{nn'}(\mathbf{k})$ , that is, find the Hamiltonian matrix  $H(\mathbf{k})_{nn'} = \langle u_{n0} | H(\mathbf{k}) | u_{n'0} \rangle$ .
4. [1 point] Calculate effective mass for 1D linear chain with one orbital in unit cell and find  $k$ -vectors within the 1<sup>st</sup> Brillouin zone where the inverse effective mass is zero. What would be the effective mass in case of isolated atoms?
5. [1 point] Show that  $f(\mathbf{r} - \mathbf{r}_0) = \exp(-\mathbf{r} \cdot \nabla)f(\mathbf{r})$ . *Hint:* use Taylor expansion of the  $f(\mathbf{r} - \mathbf{r}_0)$  around  $\mathbf{r}_0$ .
6. [2 points] Suppose 1D crystal with a lattice constant  $a$  and dispersion relation 
$$\epsilon(k) = \frac{\hbar^2}{ma^2} \left[ \frac{7}{8} - \cos(ka) + \frac{1}{8} \cos(2ka) \right].$$
  - a) Determine effective masses at the bottom and top of the band.
  - b) Express Bloch velocities at the bottom and top of the band in terms of effective mass.
7. [1 point] From definition of the Wannier functions  $\phi_n(\mathbf{r} - \mathbf{R}_j) = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{-i\mathbf{k} \cdot \mathbf{R}_j} \psi_{n\mathbf{k}}(\mathbf{r})$ , where  $n$  is the band index and  $\psi_{n\mathbf{k}}(\mathbf{r})$  is the state describing fulfilling the Bloch theorem, prove the orthogonality  $\langle \phi_n(\mathbf{r} - \mathbf{R}_j) | \phi_{n'}(\mathbf{r} - \mathbf{R}_{j'}) \rangle = \delta_{nn'} \delta_{jj'}$ .
8. [1 point] The original normalized Bloch state can be reconstructed in the form 
$$\psi_{n\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{N}} \sum_j e^{i\mathbf{k} \cdot \mathbf{R}_j} \phi_n(\mathbf{r} - \mathbf{R}_j),$$
 show that dispersion  $\epsilon(k) = \epsilon_0 + \sum_{\mathbf{R}} t(\mathbf{R})e^{-i\mathbf{k} \cdot \mathbf{R}}$ , and define  $\epsilon_0$  and  $t(\mathbf{R})$ .
9. [5 extra points] Low energy electrons in graphene are called Dirac or massless electrons and are described by the linear dispersion in momentum  $\epsilon(\mathbf{k}) = v_F|\mathbf{k}|$ , where  $v_F$  is the Fermi velocity, a material constant. Find an alternative definition of the effective mass to cure the diverging problem when using the standard definition  $(1/m)^* = 1/\hbar^2(\partial^2\epsilon(k)/\partial k^2)$ .

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